

# Quantum Abacus

Qubit Realization with  
Quantum Point Interactions

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# Introduction

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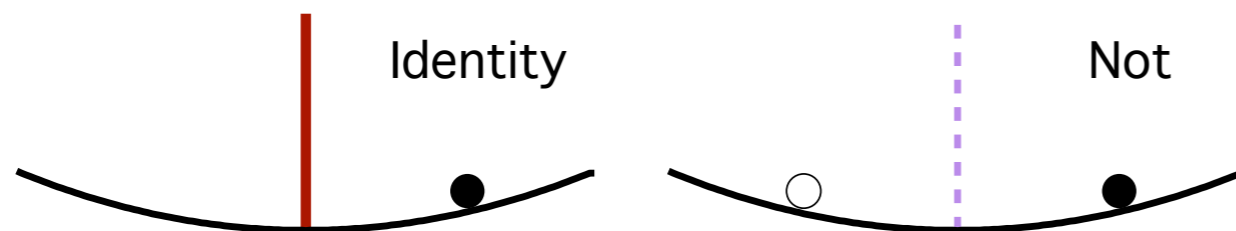
- Theoretical design of semiclassically robust qubit
- Self-driven classical **abacus** based on “locational” bit controlled by gate
- Quantum abacus based on recurring quantum states, control by **quantum gate** that is described by  $U(2)$  point interaction

# Classical abacus

- hard ball in harmonic oscillator stores a BIT as right/left location



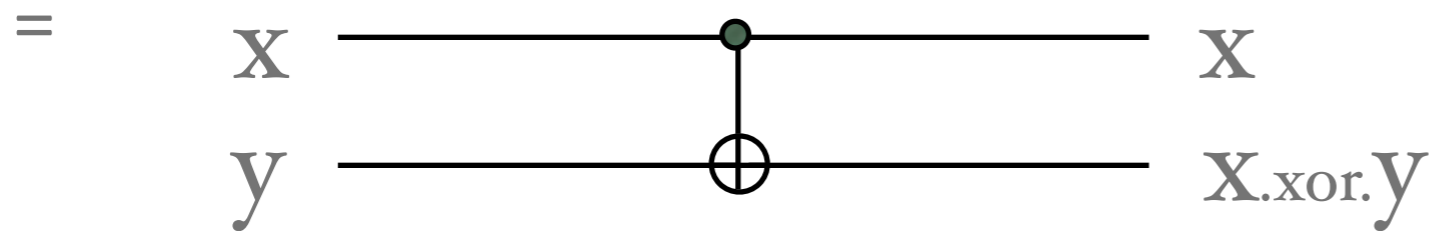
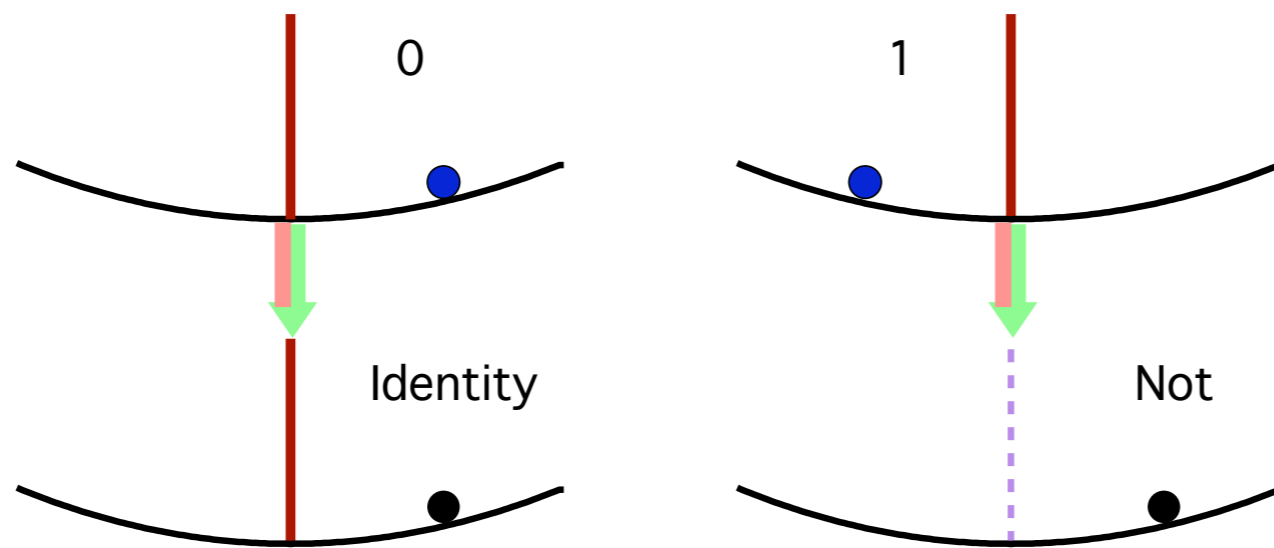
- opening/closing of hard wall GATE as logic operation in time-step  $\tau = \pi/\omega$



half oscillation period

# Two-bit trigger

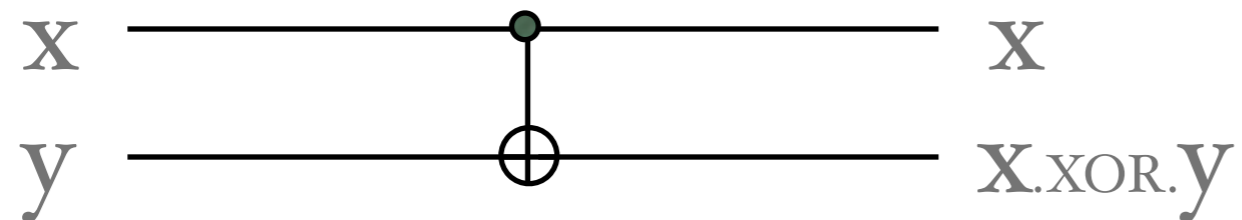
- A two-bit trigger for C-Not operation



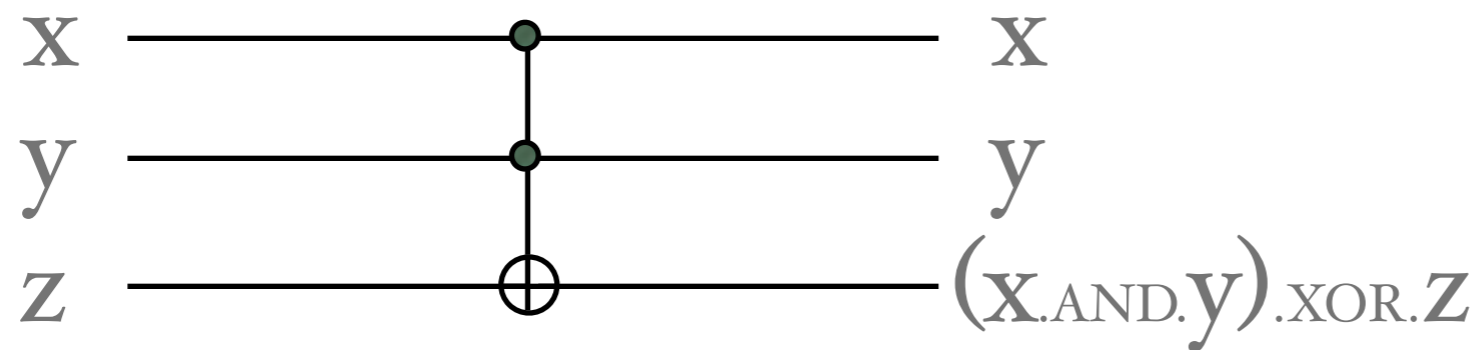
# Three-bit trigger

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- In a similar manner to C<sub>N</sub>T<sub>L</sub>-Not gate

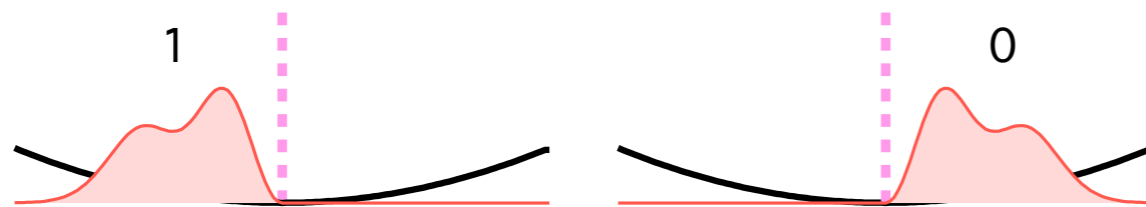


three-bit trigger gate can be designed to create universal C<sub>N</sub>T<sub>L</sub>-C<sub>N</sub>T<sub>L</sub>-Not gate

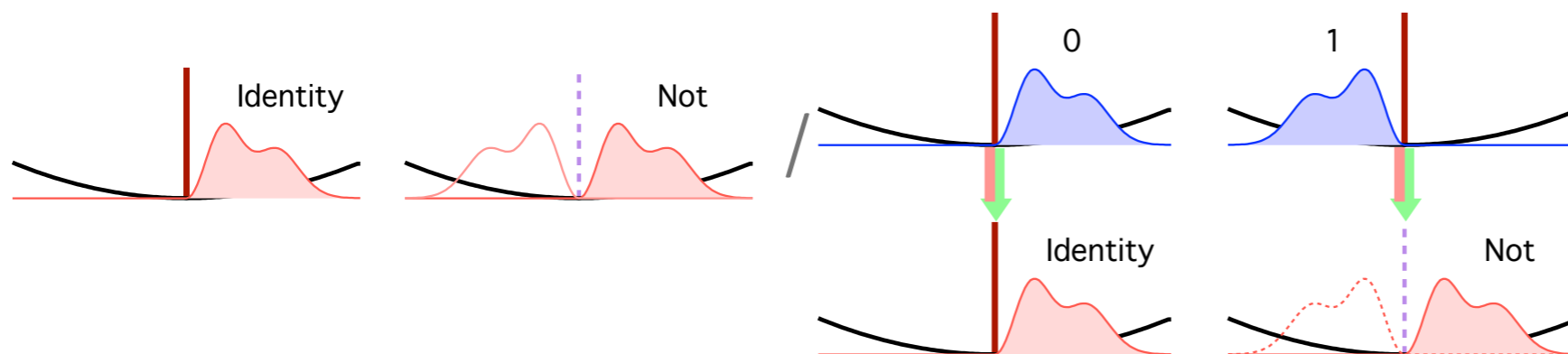


# Quantum waves

- works just fine if localized in one side



... with same gate and trigger

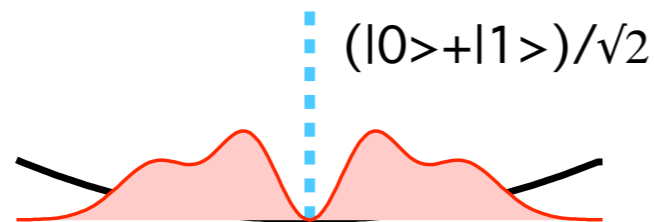


- ... since states get mirrored after  $\tau$ :  
 $\psi(x, \tau) = \pm \psi(x, 0)$  with h.o. potential

# Qubit superposition

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- In general, however, a quantum state is bilocalized



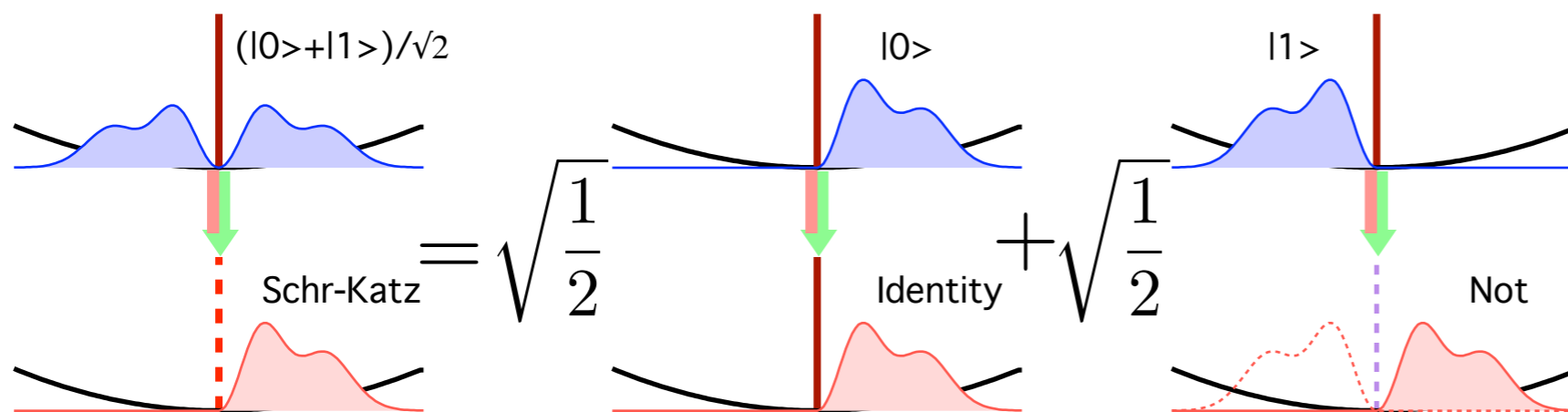
describable as sum of **QUBITS**  $|0\rangle$  and  $|1\rangle$

$$= \sqrt{\frac{1}{2}} \text{ (diagram for } |0\rangle \text{)} + \sqrt{\frac{1}{2}} \text{ (diagram for } |1\rangle \text{)}$$

- An advantage might be taken
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# Schrödinger-cat gate

- If, instead, bilocal wave is used for trigger gate



...is conditionally half open/closed gate

- parallel multiple operation

# U(2) Hilbert space

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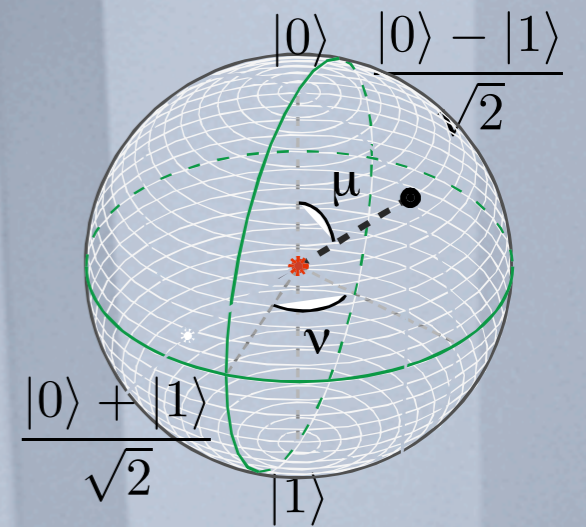
- Quantum information processing possible, if we realize/control U(2) states

$$|\Psi\rangle = \alpha |0\rangle + \beta |1\rangle$$

- Essentially, apart from overall phase, need to realize **Bloch sphere**  $\{\mu, \nu\}$

$$|\Psi\rangle = \cos \frac{\mu}{2} |0\rangle + e^{i\nu} \sin \frac{\mu}{2} |1\rangle$$

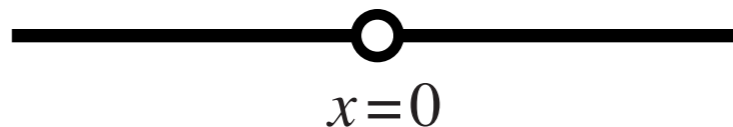
- This is achieved by .....
- 



# Quantum barriers

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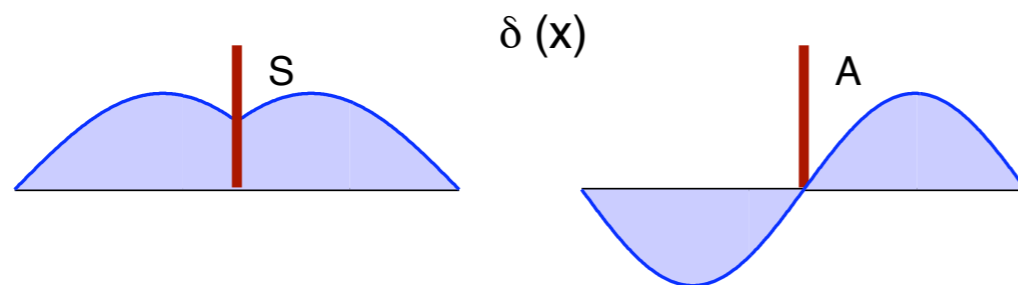
- Regulate the temporal development of wave function at each time-step  $\tau$  by controlling the **characteristics** of ...
- Quantum mechanically possible **thin barriers**, which is idealized as ...
- 1D **Quantum Point Interactions** ...



- that generalize open/closed gate

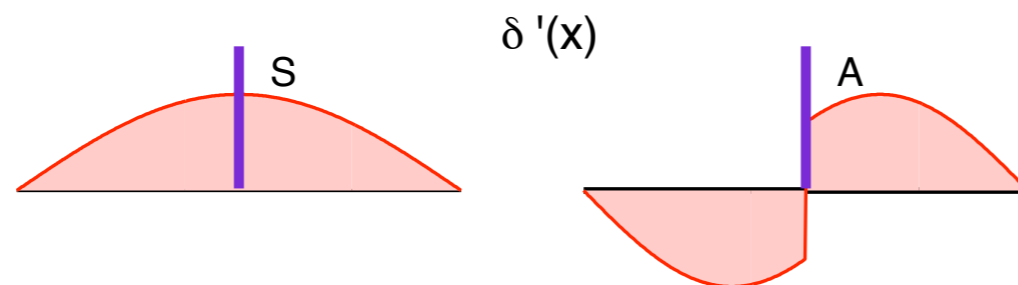
# Some point interactions

- Delta-interaction :  $\Gamma$ -parameter family



Discontinuity in  $\psi'(x)$  » Dirichlet wall

- Delta'-interaction :  $\Gamma$ -parameter family



Discontinuity in  $\psi(x)$  » Neumann wall

# Matrix notations

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- Express qubits as 2-vectors

$$|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad |1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

- Some Bloch-sphere elements

$$I = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} : \text{Identity} \quad \sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} : \text{Not}$$

$$h = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} : \text{Hadamard}$$

# U(2) point interactions

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- Right/left vectors

$$\Psi(x) = \begin{pmatrix} \psi(x) \\ \psi(-x) \end{pmatrix}, \Psi'(x) = \begin{pmatrix} \psi'(x) \\ -\psi'(-x) \end{pmatrix}, \quad (x > 0)$$

- **Current conservation** at  $x=0$  gives the **connection condition**

$$(U - I)\Psi(0) + iL_0(U + I)\Psi'(0) = 0$$

- $U \in U(2) \approx S^1 \times S^3$  : 4-parameter family

arbitrary scale parameter

# SU(2) algebra

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- Consider transformations on  $\psi(x)$ 
  - $P_1: \psi(x) \rightarrow \psi(-x)$  : mirror-reflection
  - $P_2: \psi(x) \rightarrow -i(\Theta(x)-\Theta(-x))\psi(-x)$
  - $P_3: \psi(x) \rightarrow (\Theta(x)-\Theta(-x))\psi(x)$  : half-flip
- They act on  $\Psi$  as Pauli matrices
  - $P_1: \Psi(x) \rightarrow \sigma_1 \Psi(x)$
  - $P_2: \Psi(x) \rightarrow \sigma_2 \Psi(x)$
  - $P_3: \Psi(x) \rightarrow \sigma_3 \Psi(x)$

# Spectral invariance

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- Let  $H_U$  be Hamiltonian with  $U$  and let  $\sigma$  be one of  $\sigma_i$
  - $U' = \sigma U \sigma \iff H_{U'} = P H_U P$
  - $H_U$  and  $H_{U'}$  share the **same spectra** because  $H_{U'} P = P H_U$   
 $H_U \psi = E \psi \iff H_{U'} (P \psi) = E (P \psi)$
  - Spectra invariant for all “spin rotation”
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# Spectral Decomposition

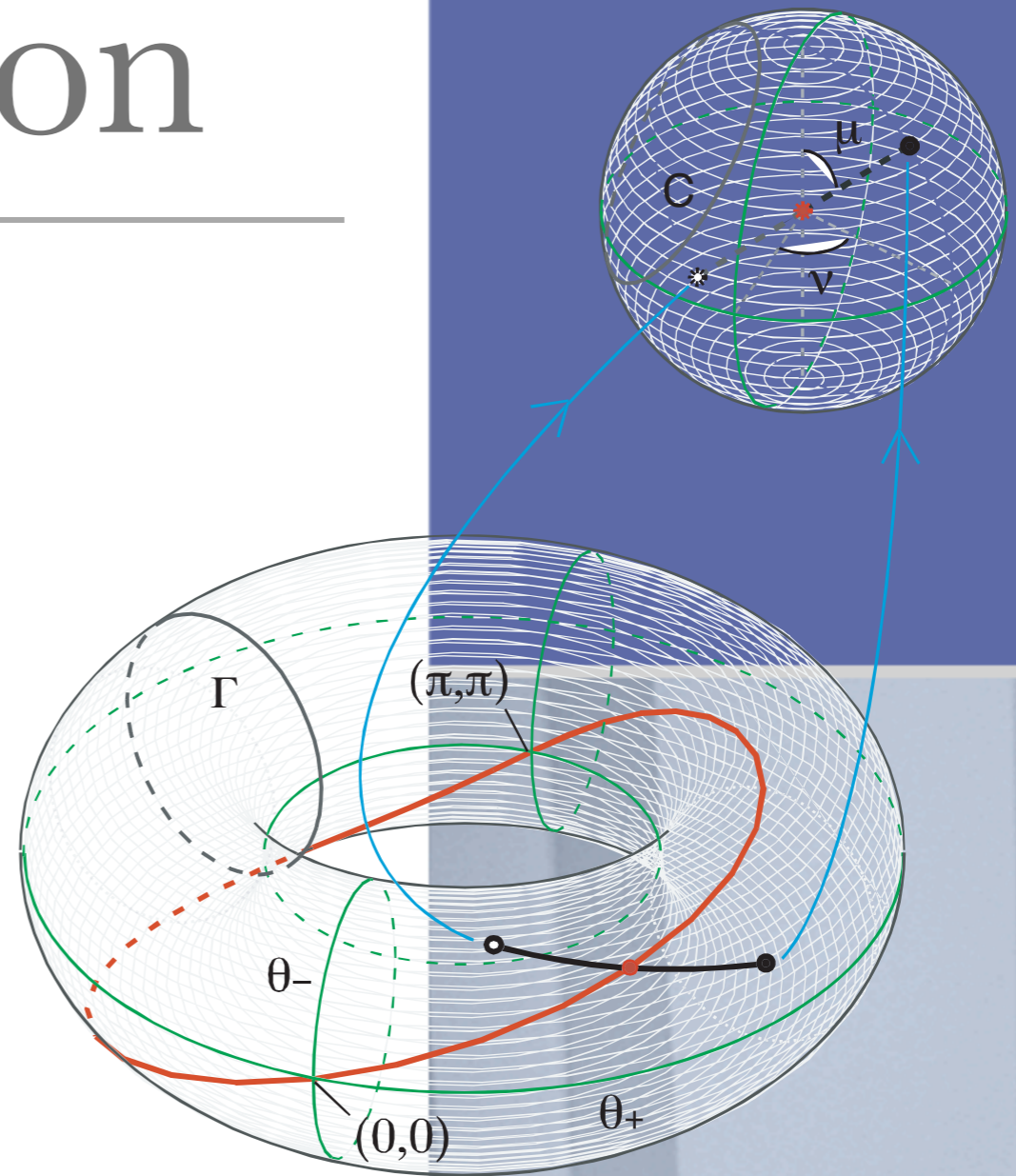
- Any element  $U$  is expressed as

$$U = \Sigma D \Sigma$$

$$\Sigma = \begin{pmatrix} \cos \frac{\mu}{2} & e^{i\nu} \sin \frac{\mu}{2} \\ e^{-i\nu} \sin \frac{\mu}{2} & -\cos \frac{\mu}{2} \end{pmatrix}$$

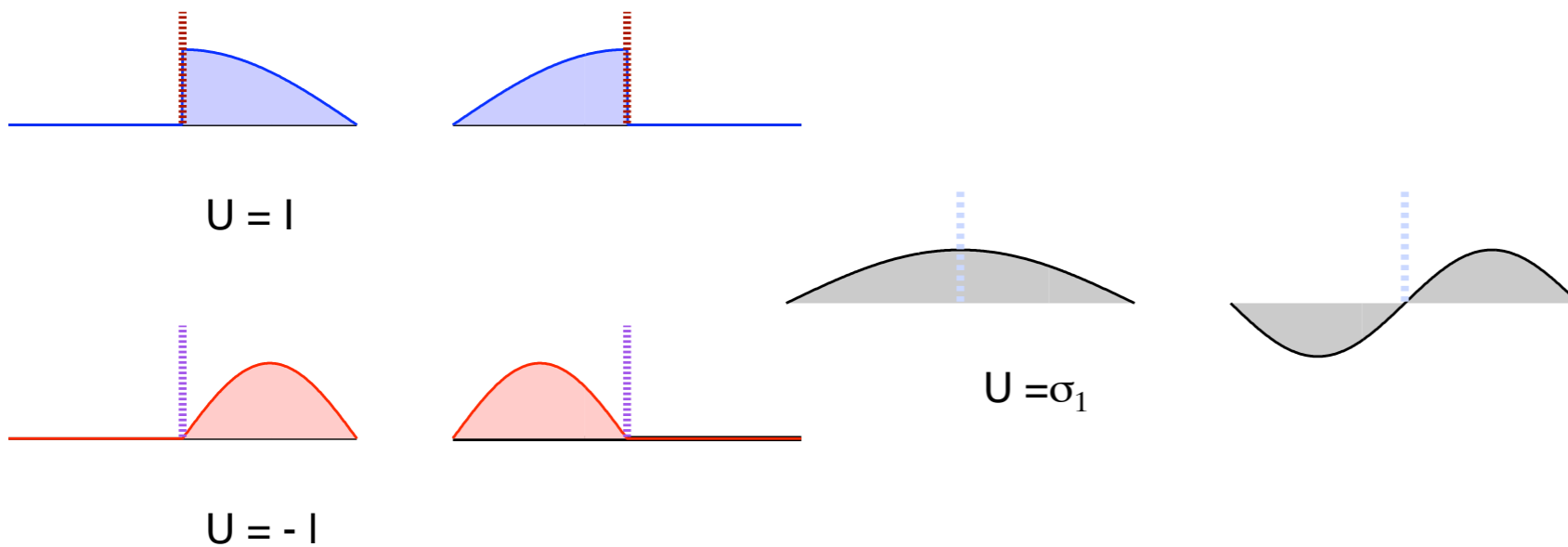
$$D = \begin{pmatrix} e^{i\theta_+} & 0 \\ 0 & e^{i\theta_-} \end{pmatrix}$$

$\{\mu, \nu\} \in S^2$  : isospectral,  $\{\theta_+, \theta_-\} \in T^2$  : spectral



# Quantum hard-barriers

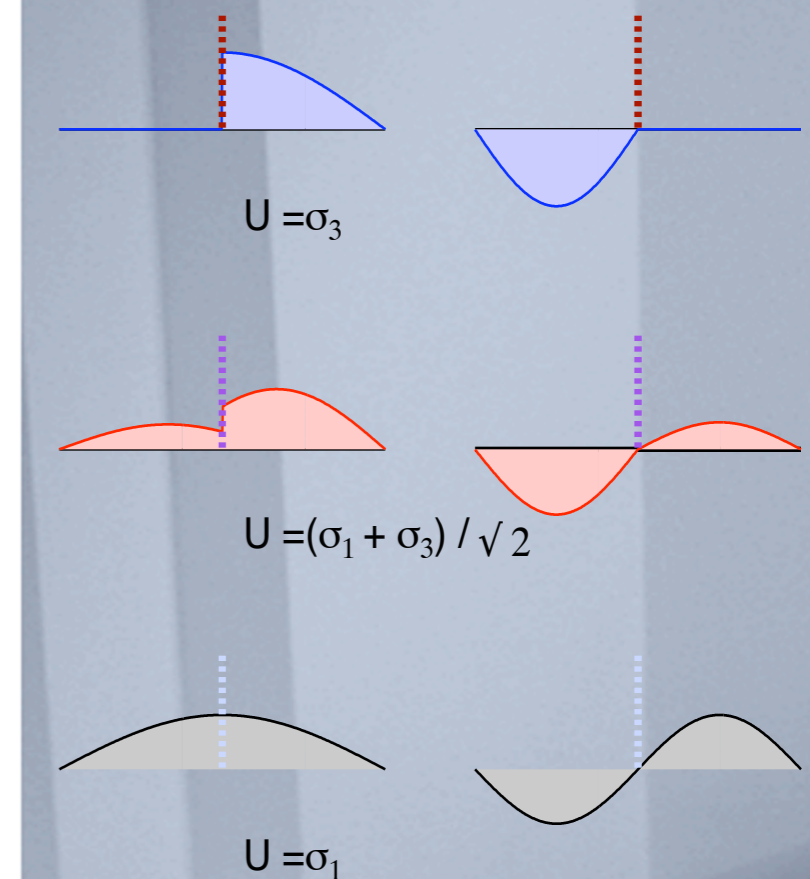
- Neumann and Dirichlet walls corresponds to  $U = I$  and  $U = -I$



- No-barrier corresponds to  $U = \sigma_I$

# Exotic contact forces

- $U = \sigma_3 \cos \frac{\mu}{2} + \sigma_1 \sin \frac{\mu}{2} \cos \nu + i\sigma_2 \sin \frac{\mu}{2} \sin \nu$   
 $U \in S^2$  (isospectral)
- Isospectral sphere is “scale-invariant”  
»» E-indep. constant transmission
- Interpolate between  $\sigma_1$  and  $\sigma_3$
- Treat them as gates to obtain desired mixture of localized wave functions



# Realizing contact forces

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- Delta interaction is a short-range limit of potential with constant volume integral
  - How can you produce other exotic point interactions?
  - Ans: by singular short-range limit of up to five delta potentials
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# Examples of realization

- Exitic delta'

$$U = I$$

- Dirac delta

$$U = -I$$

- No barrier (free)

$$U = \sigma_1$$

- Scale invariant

$$U = H$$

T.Cheon, T.Shigehara,  
Phys. Lett. A243 (1998) III.  
&  
P.Exner, H.Neidhardt, V.Zagrebnov,  
Comm. Math. Phys. 224 (2001) 593.

# State evolution

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- Time evolution of  $\Psi(x)$  under given point interaction  $U \in S^2$  (isospectral)

$$\Psi \xrightarrow[t]{} \mathcal{U}_t(U) \Psi$$

- With the time step  $\tau = \pi/\omega$

Fact 1 :  $\mathcal{U}_\tau(I) = -i$

Fact 2 :  $\mathcal{U}_\tau(\sigma_1) = -i\sigma_1$

Fact 3 :  $\mathcal{U}_\tau(\sigma_3) = -i\sigma_3$

- Parametric  $S^2$  turned into Bloch  $S^2$  ?!
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# Bloch-sphere imprinting

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- Suppose  $U' = VUV^{-1}$

Time evolution under  $U$  and  $U'$  are related by  $\mathcal{U}_\tau(U') = V\mathcal{U}_\tau(U)V^{-1}$

- $U = V\sigma_3V^{-1}$ ,  $V \in SU(2)$   
span  $S^2$ (scale inv)

- Therefore, for  $U \in S^2(\text{Bloch})$

$$\mathcal{U}_\tau(U) = -iU$$

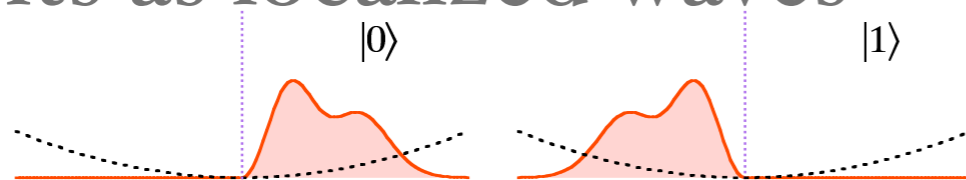
»» Qubit control by quantum gate

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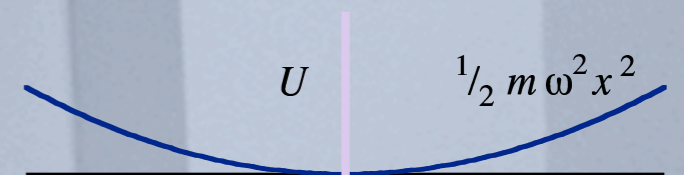
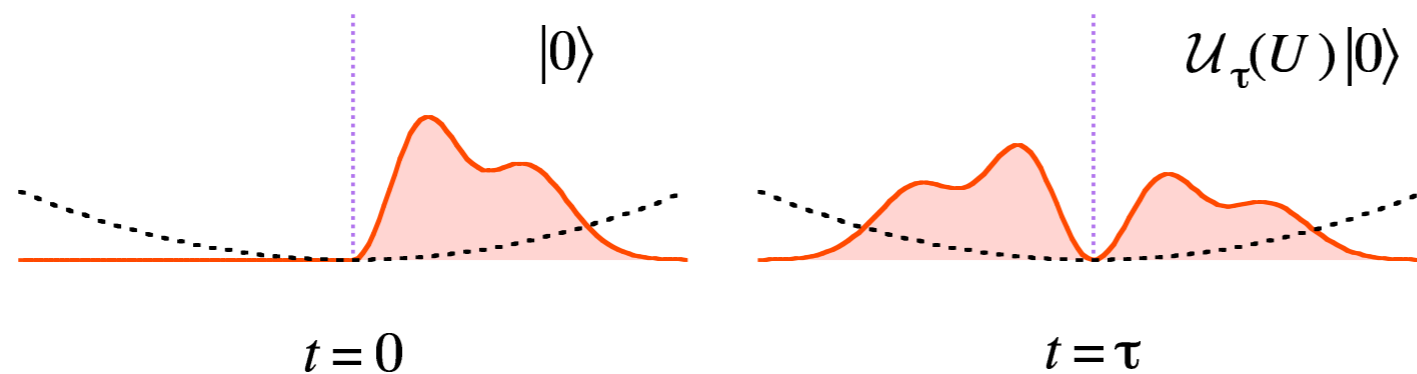
# Quantum abacus

- Harmonic oscillator and  $U$

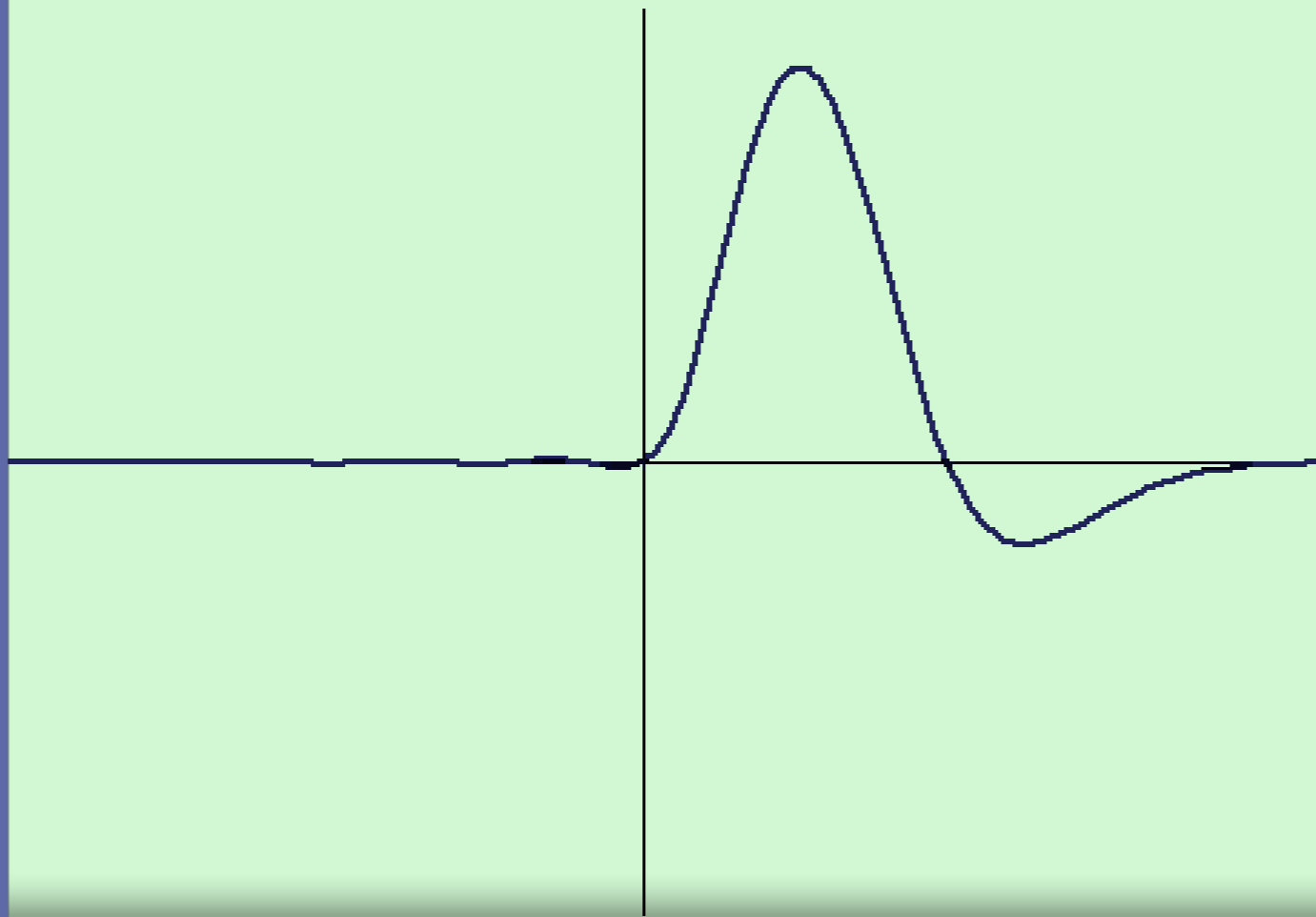
Qubits as localized waves



- Evolve under  $U$  for period  $t$  to get  $-iU$



$U = H$  (Hadamard)

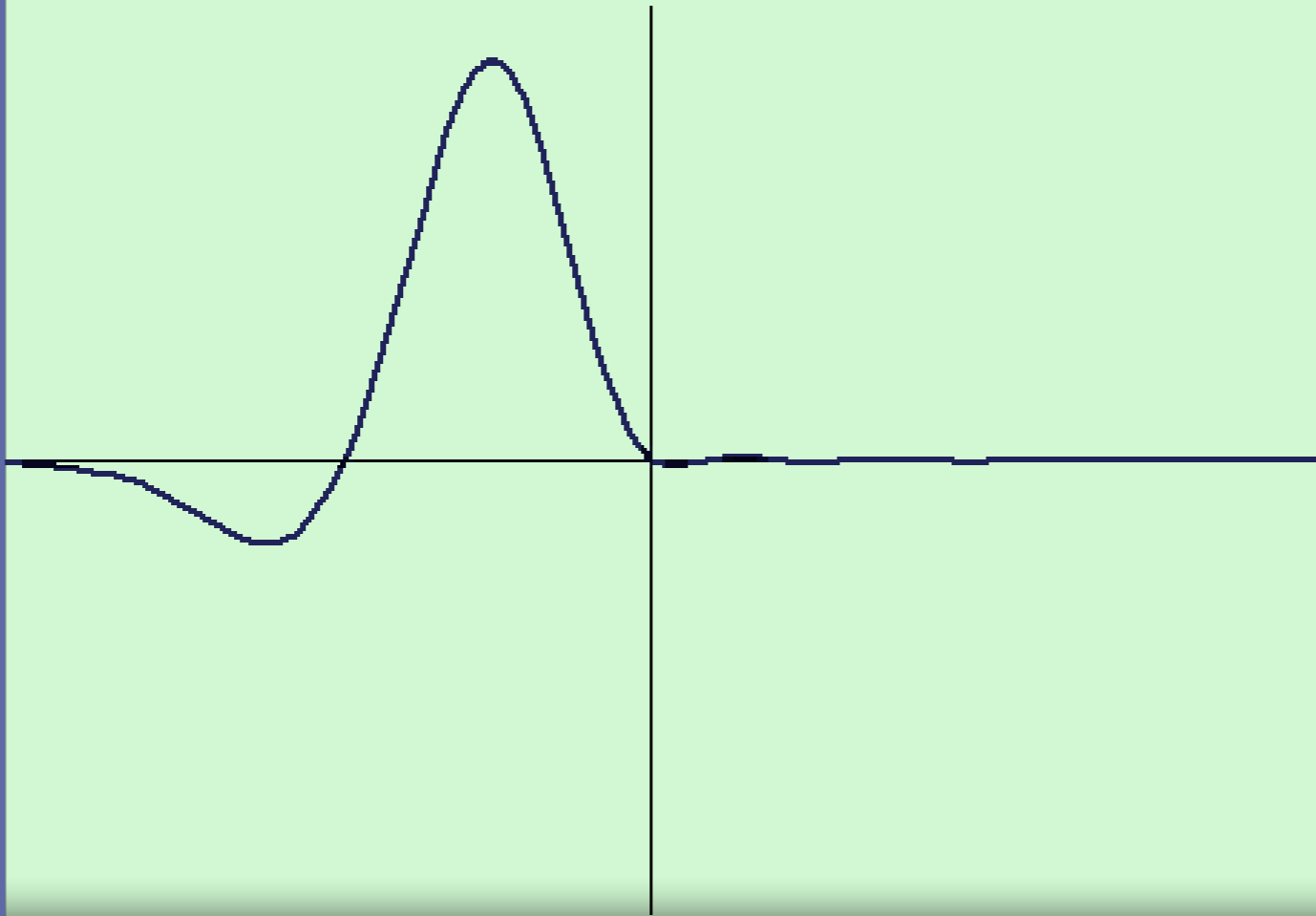


Quantum Operations  $U|0\rangle$

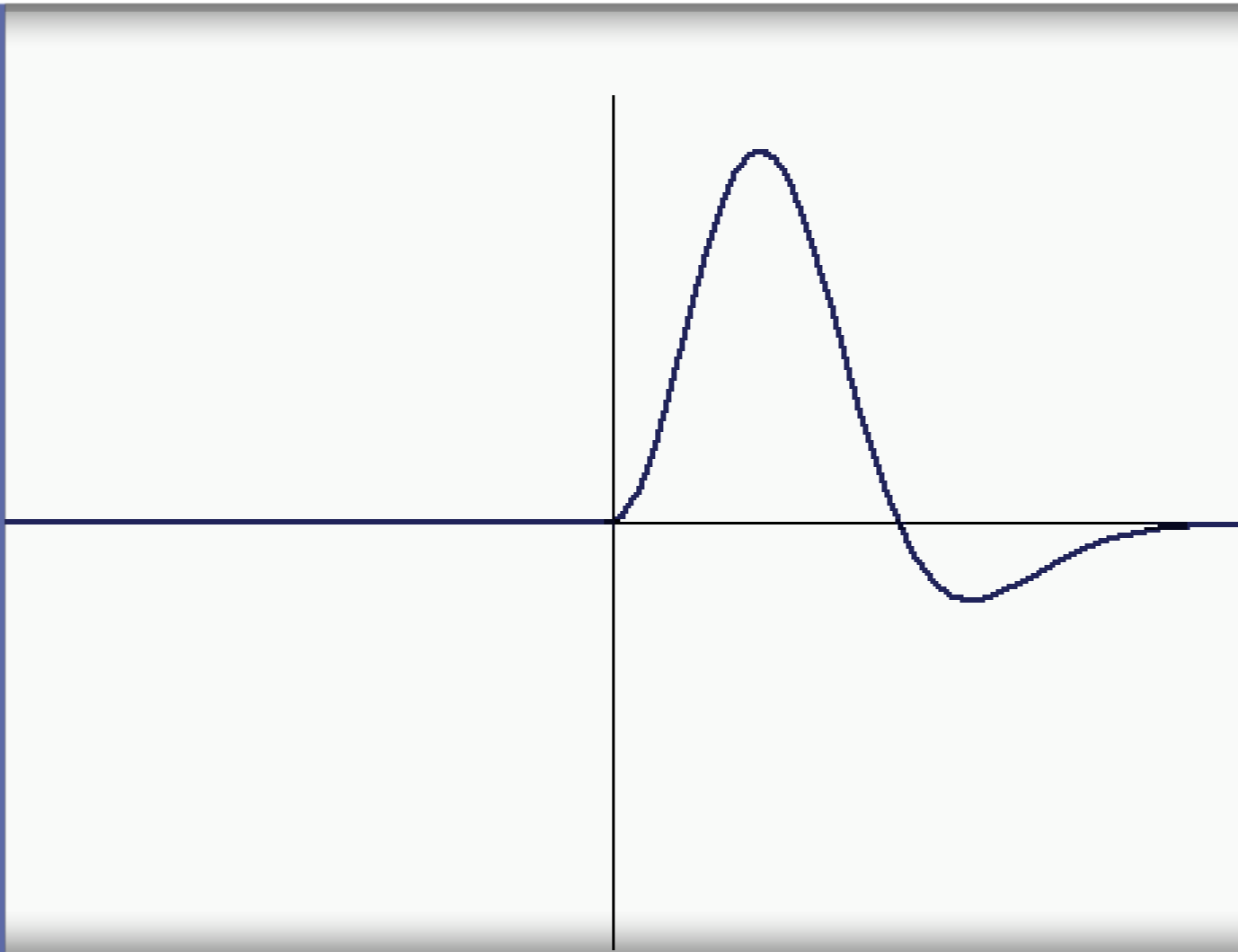
Movie:

24

$U = H$  (Hadamard)



Quantum Operations  $U|1\rangle$   
Movie:  
25



Movie:  
Abacus Operation .. $U_3U_2U_1 | 0 \rangle$   
26

# Summary

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- Quantum qubit is constructed on 1D line with harmonic potential which is divided by a barrier, or ...
  - Quantum point interaction, whose mathematical structure is  $U(2)$ .
  - Parameter space  $S^2$  of scale invariant barriers is turned into Hilbert space sphere through Bloch imprinting
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# References

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- T. Cheon, T. Fülöp and I. Tsutsui, “Quantum abacus”, arXiv.org, quant-ph/0404039 (2004).
- T. Cheon, T. Fülöp and I. Tsutsui, “Symmetry, duality and anholonomy of point interaction in one-dimension”, Ann of Phys. (NY) **294** (2001) 1-23.
- I. Tsutsui, T. Fülöp and T. Cheon, “Möbius structure the spectral space of Schrödinger operator with of point interaction”, J. Math. Phys. **42** (2001) 5687-5697.
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